

Hadrons and Holography

Josh Erlich

College of William & Mary

MENU 2010, Williamsburg VA

What is Holographic QCD?

- Holographic QCD is an attempt to model hadronic physics as a theory of fields or strings in extra dimension(s).

What is Holographic QCD?

- Holographic QCD is an attempt to model hadronic physics as a theory of fields or strings in extra dimension(s).
- Simple models of this type capture qualitative, and sometimes quantitative, features of QCD at low energies.

QCD

Holographic QCD

Towers of bound states identified
by quantum numbers, mass

Towers of Kaluza-Klein modes
identified by quantum numbers, mass

QCD

Holographic QCD

Towers of bound states identified by quantum numbers, mass	Towers of Kaluza-Klein modes identified by quantum numbers, mass
Hidden local symmetry: Vector mesons act like massive gauge bosons (Sakurai; Bando et al.)	Vector mesons modeled by Kaluza-Klein modes of gauge fields (Polchinski, Strassler; Son, Stephanov; Brodsky, De Teramond; etc.)

QCD

Holographic QCD

Towers of bound states identified by quantum numbers, mass	Towers of Kaluza-Klein modes identified by quantum numbers, mass
Hidden local symmetry: Vector mesons act like massive gauge bosons (Sakurai; Bando et al.)	Vector mesons modeled by Kaluza-Klein modes of gauge fields (Polchinski, Strassler; Son, Stephanov; Brodsky, De Teramond; etc.)
Chiral symmetry breaking has implications at low energies	Chiral symmetry breaking has implications at low energies

QCD

Holographic QCD

Towers of bound states identified by quantum numbers, mass	Towers of Kaluza-Klein modes identified by quantum numbers, mass
Hidden local symmetry: Vector mesons act like massive gauge bosons (Sakurai; Bando et al.)	Vector mesons modeled by Kaluza-Klein modes of gauge fields (Polchinski, Strassler; Son, Stephanov; Brodsky, De Teramond; etc.)
Chiral symmetry breaking has implications at low energies	Chiral symmetry breaking has implications at low energies
Weinberg sum rules	Weinberg sum rules

QCD

Holographic QCD

Towers of bound states identified



Weinberg sum rules

Towers of Kaluza-Klein modes



Weinberg sum rules

QCD

Holographic QCD

Towers of bound states identified



Towers of Kaluza-Klein modes



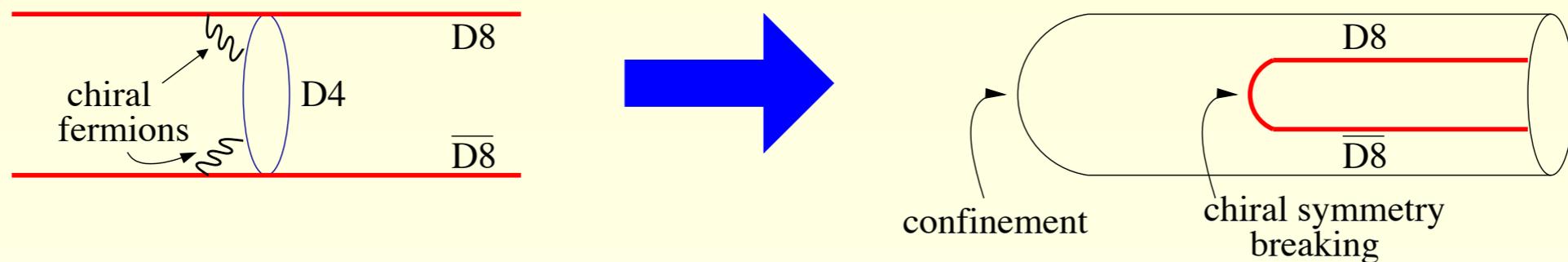
Turkey

Tofu

Top-Down AdS/QCD

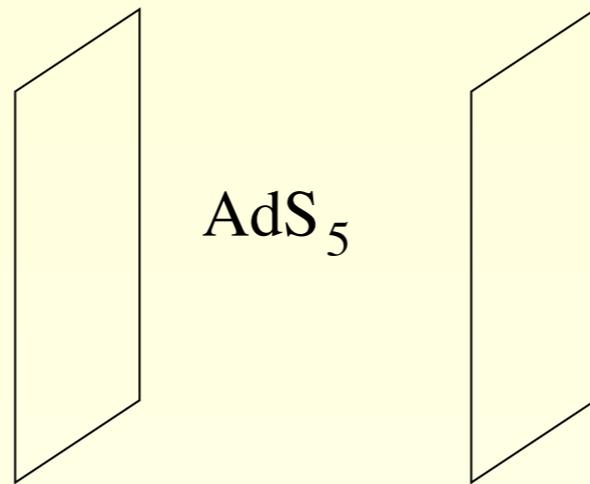
- String theory brane configuration \rightarrow gauge theory similar to QCD (e.g. [Kruczenski et al.](#); [Antonyan, Harvey, Kutasov](#); [Sakai, Sugimoto](#))
- At **large- N** , theory has weakly-coupled dual description via the AdS/CFT correspondence ([Maldacena](#))

The Sakai-Sugimoto Model



Bottom-Up AdS/QCD

- Model tower of resonances as Kaluza-Klein modes in an extra dimension (Son,Stephanov'04)
- Model pattern of chiral symmetry breaking by analogy with AdS/CFT correspondence
- *Optional*: Specify details of model (geometry of extra dimension, couplings) by matching to UV as best possible (e.g. Brodsky,De Teramond; JE *et al.*; Da Rold,Pomarol)



Top-Down AdS/QCD:

- Advantage: Both descriptions of theory are relatively well understood, duality is exact.
- Disadvantage: QCD with fundamental flavors does not have weakly-coupled AdS/CFT dual, so far even at large- N .

Bottom-Up AdS/QCD:

- Advantage: Freedom to match model to aspects of QCD.
- Disadvantage: Some features of model disagree with QCD (analogous to large- N limit).

Building a Bottom-Up Model

Step 1: Choose 5D gauge group and geometry.

- Tower of vector mesons are identified with tower of Kaluza-Klein gauge bosons.

SU(2) isospin \rightarrow 5D SU(2) gauge theory

Conformal in the UV \rightarrow Anti-de Sitter space near its boundary

Building a Bottom-Up Model

Step 1: Choose 5D gauge group and geometry.

- Tower of vector mesons are identified with tower of Kaluza-Klein gauge bosons.

SU(2) isospin \rightarrow 5D SU(2) gauge theory

Conformal in the UV \rightarrow Anti-de Sitter space near its boundary

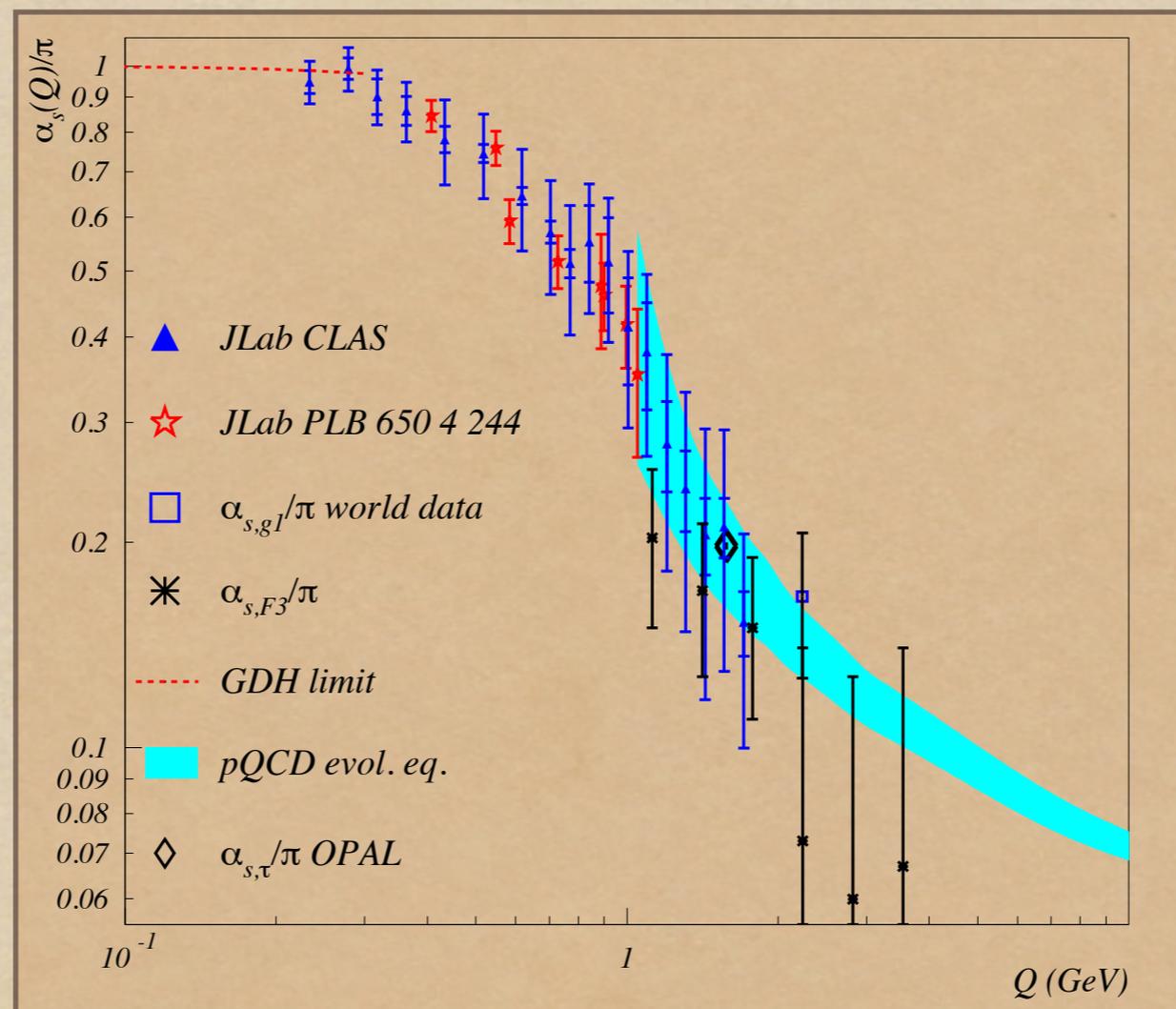
Can choose geometry by matching spectrum to Pade approx of SU(2) current-current correlator in deep Euclidean regime $-q^2 \gg m_\rho^2$.

Result: geometry = slice of AdS space

(Shifman; JE, Kribs, Low; Falkowski, Perez-Victoria).

Evidence for Conformality

Brodsky and collaborators motivate Anti-de Sitter space from approximate conformality of QCD at low energies. e.g. Brodsky and Shrock '08



From CLAS (Deur et al.) '08

Building a Bottom-Up Model

To include the full chiral symmetry, not just the vector subgroup,

$SU(2) \times SU(2)$ chiral symmetry \rightarrow $SU(2) \times SU(2)$ 5D gauge group

Additional tower of gauge bosons \rightarrow tower of axial-vector mesons.
(5D parity \rightarrow 4D parity)

(Also describes pions after symmetry breaking)

Building a Bottom-Up Model

Step 2: Include pattern of chiral symmetry breaking

Hint from AdS/CFT: 4D operator \rightarrow 5D field

$\bar{q}_i q_j \rightarrow$ Scalar fields X_{ij} , bifundamental under $SU(2) \times SU(2)$

Background profile for X_{ij} :

Non-normalizable mode \rightarrow source $\mathcal{L}_{4D} \supset m_{ij} \bar{q}_i q_j$

Normalizable mode \rightarrow VEV $\langle \bar{q}_i q_j \rangle$

The scalar field background explicitly and spontaneously breaks the chiral symmetry.

Building a Bottom-Up Model

For definiteness, we need to choose 5D mass of scalar field.

AdS/CFT: $m_X^2 = \Delta_{\bar{q}q}(\Delta_{\bar{q}q} - 4)$ in units of AdS curvature.

In the UV, $\Delta_{\bar{q}q} = 3$, so we choose $m_X^2 = -3$.

Note: This choice is made for definiteness, but is not necessary.

Building a Bottom-Up Model

In summary, the model is:

$SU(2) \times SU(2)$ gauge theory in slice of AdS_5 with background bifundamental scalar field.

$$S = \int d^5x \sqrt{-g} \left(-\frac{1}{2g_5^2} \text{Tr} (L_{MN}L^{MN} + R_{MN}R^{MN}) + \text{Tr}(|D_M X|^2 - 3|X|^2) \right)$$

$$ds^2 = \frac{1}{z^2} (dx_\mu dx^\mu - dz^2), \quad \epsilon < z < z_{IR}$$

$$X_0(x, z) = \frac{m_q}{2} z + \frac{\langle \bar{q}q \rangle}{2} z^3$$

Model parameters: $g_5, m_q, \langle \bar{q}q \rangle, z_{IR}$

(JE, Katz, Son, Stephanov; DaRold, Pomarol)

Building a Bottom-Up Model

AdS/QCD reproduces consequences of chiral symmetry, e.g.
Gell-Mann, Oakes, Renner relation

$$m_{\pi}^2 f_{\pi}^2 = 2m_q \langle \bar{q}q \rangle$$

Matching to UV

In the deep Euclidean regime $-q^2 \gg m_\rho^2$, perturbative QCD gives

$$i \int d^4x e^{iq \cdot x} \langle J_\mu^a(x) J_\nu^b(0) \rangle = (q_\mu q_\nu - g_{\mu\nu} q^2) \delta^{ab} \frac{N}{24\pi^2} \log(q^2)$$

We can express the correlator as a sum over resonances:

$$i \int d^4x e^{iq \cdot x} \langle J_\mu^a(x) J_\nu^b(0) \rangle = \sum \frac{F_n^2}{q^2 - m_n^2} \left(g_{\mu\nu} - \frac{q_\mu q_\nu}{m_n^2} \right) \delta^{ab}$$

Agreement of these expressions in the deep Euclidean regime is a Weinberg sum rule.

$m_n = n^{\text{th}}$ Kaluza-Klein mass

$F_n =$ Decay constant of n^{th} resonance

Matching to UV

Matching 5D calculation w/ 4D perturbative calculation in UV \rightarrow

$$g_5^2 = 12\pi^2/N.$$

Note: this choice is made for definiteness, but is not necessary.

Summary so far:

We have constructed a three-parameter model of the ρ , a_1 and π

Soft-Wall AdS/QCD

In the Hard Wall model $m_n^2 \sim n^2$

To obtain a linear Regge trajectory, the geometry can be modified while coupling to a dilaton background.

(Karch, Katz, Son, Stephanov '06)

$$S = \int d^5x \sqrt{g} e^{-\Phi(x,z)} \mathcal{L}$$

$$\Phi_0(z) \sim z^2, \quad g_{MN} = \text{AdS}_5 \text{ Metric}$$

Low-energy predictions are similar to hard-wall model

Hard-Wall (5D tree level)

With $z_{IR} = 1/(346 \text{ MeV})$, $\langle \bar{q}q \rangle = (308 \text{ MeV})^3$, $m_q = 2.3 \text{ MeV}$

Central Values

Observable	Measured (MeV)	Model (MeV)
m_π	139.6	141
m_ρ	775.8	832
m_{a_1}	1230	1220
f_π	92.4	84.0
$F_\rho^{1/2}$	345	353
$F_{a_1}^{1/2}$	433	440
$g_{\rho\pi\pi}$	6.03	5.29
m_{f_2}	1275	1236

From JE et al. '05, Katz et al. '05

Hard-Wall (5D tree level)

With strange quark mass parameter

Observable	Measured (MeV)	Model (MeV)
m_{K^*}	892	897
m_ϕ	1020	994
m_{K_1}	1272	1290
m_K	498	411
f_K	113	117
m_{ω_3}	1667	1656
m_{f_4}	2025	2058
m_η	548	520
m'_η	958	867

from E. Katz, Lattice 2008

Hard-Wall (5D tree level)

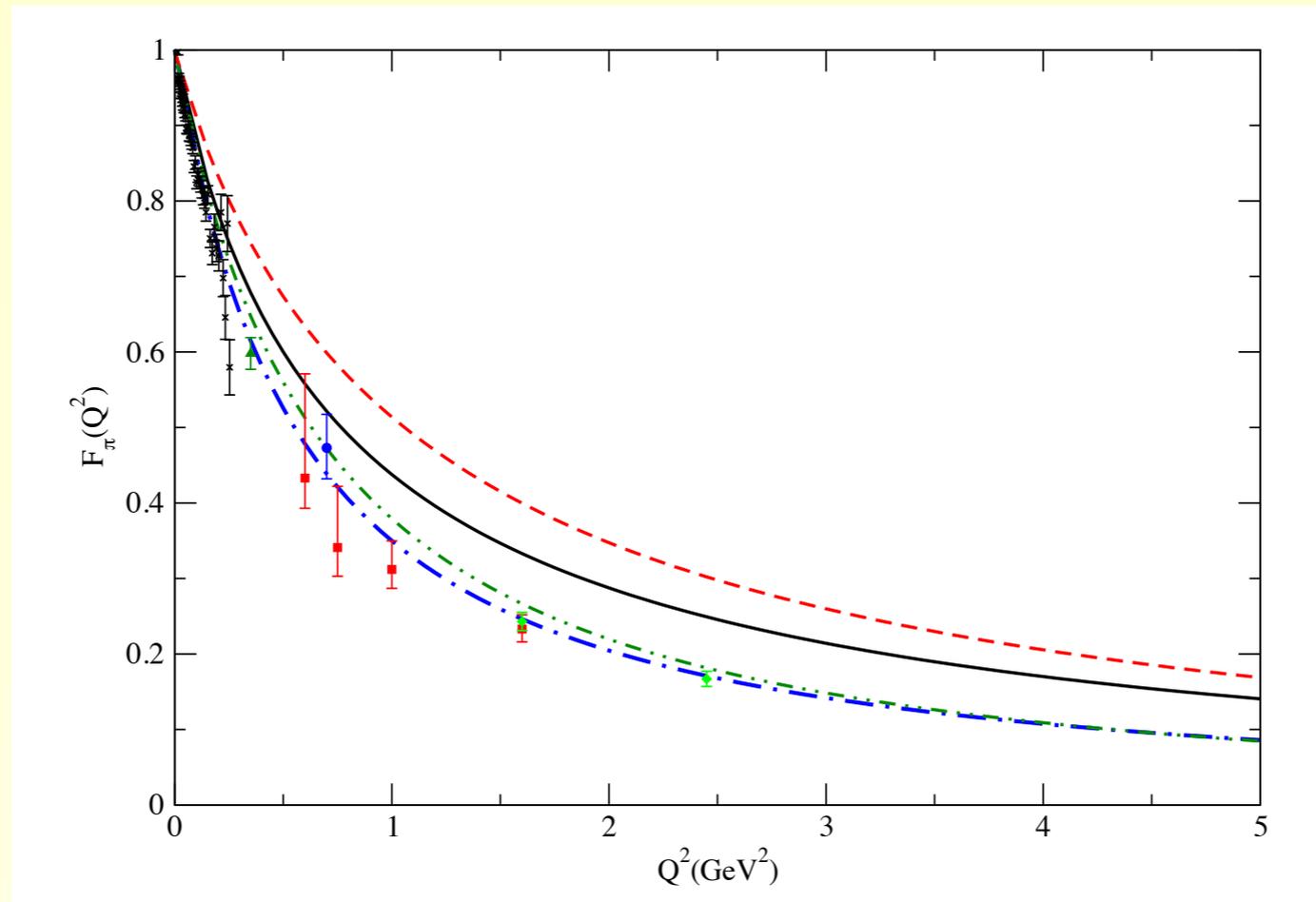
With strange quark mass parameter

Observable	Model A ($\sigma_s = \sigma_q$) (MeV)	Model B ($\sigma_s \neq \sigma_q$) (MeV)	Measured (MeV)
m_π	(fit)	134.3	139.6
f_π	(fit)	86.6	92.4
m_K	(fit)	513.8	495.7
f_K	104	101	113 ± 1.4
$m_{K_0^*}$	791	697	672
$f_{K_0^*}$	28.	36	
m_ρ	(fit)	788.8	775.5
$F_\rho^{1/2}$	329	335	345 ± 8
m_{K^*}	791	821	893.8
$F_{K^*}^{1/2}$	329	337	
m_{a_1}	1366	1267	1230 ± 40
$F_{a_1}^{1/2}$	489	453	433 ± 13
m_{K_1}	1458	1402	1272 ± 7
$F_{K_1}^{1/2}$	511	488	

From Abdidin and Carlson '09

Hard/Soft-Wall (5D tree level)

Pion Form Factor



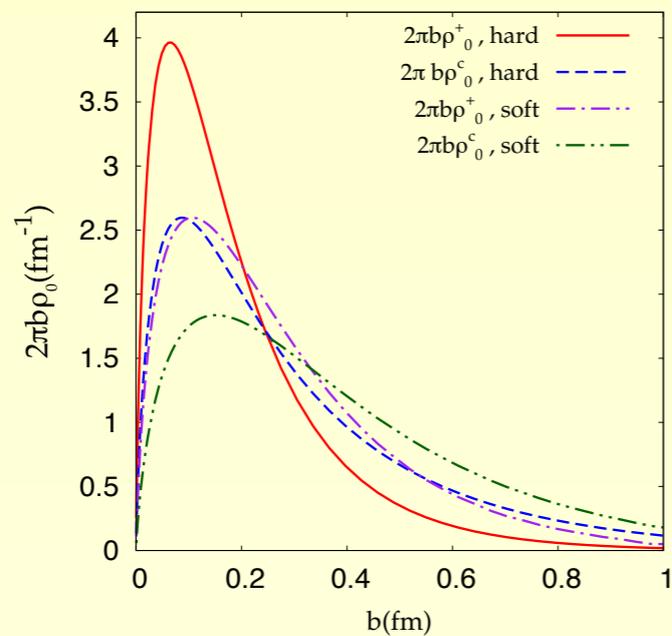
from Kwee and Lebed, arXiv:0807.4565

Solid black and blue curves: Hard wall model

Dotted red and green curves: Soft wall model

See also Grigoryan, Radyushkin '08

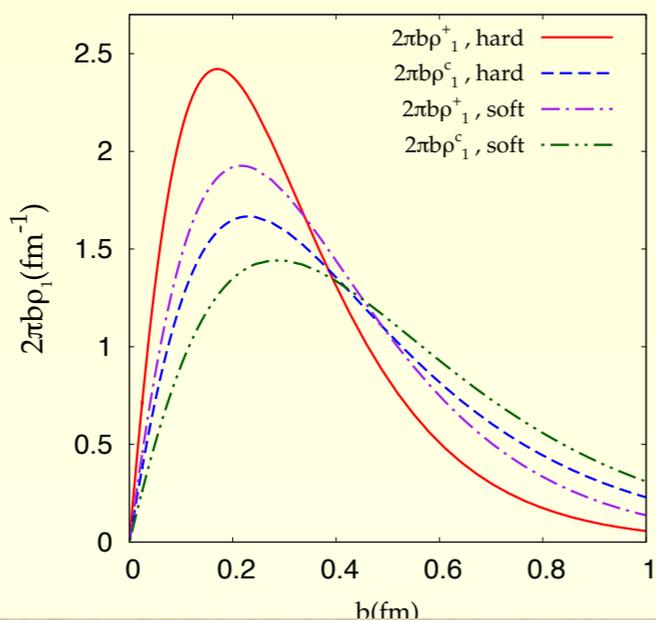
Hard/Soft-Wall (5D tree level)



Gravitational Form Factors and Generalized Parton Distributions

From Abidin and Carlson,
arXiv:0801.3839

Top: p^+ and charge densities of Helicity-0 rho mesons in hard and soft wall models



Bottom: Same for Helicity-1 rho mesons

Hard-Wall (5D tree level)

Can determine meson **radii** from behavior of form factors near $q^2 = 0$.

Hard wall model:

$$\langle r_\pi^2 \rangle_{charge} = 0.33 \text{ fm}^2$$

$$\langle r_\pi^2 \rangle_{grav} = 0.13 \text{ fm}^2$$

$$\langle r_\rho^2 \rangle_{charge} = 0.53 \text{ fm}^2$$

$$\langle r_\rho^2 \rangle_{grav} = 0.21 \text{ fm}^2$$

$$\langle r_{a_1}^2 \rangle_{charge} = 0.39 \text{ fm}^2$$

$$\langle r_{a_1}^2 \rangle_{grav} = 0.15 \text{ fm}^2$$

H. Grigoryan and A. Radyushkin; Z. Abidin and C. Carlson '07,'08

Universality in AdS/QCD?

Some observables are truly universal, *i.e.* independent of details of model.

Famous Example: Viscosity to Entropy Density η/s

Finite temperature \rightarrow spacetime horizon

Prediction, independent of details of spacetime geometry:

$$\frac{\eta}{s} = \frac{1}{4\pi}$$

Kovtun, Son and Starinets '04

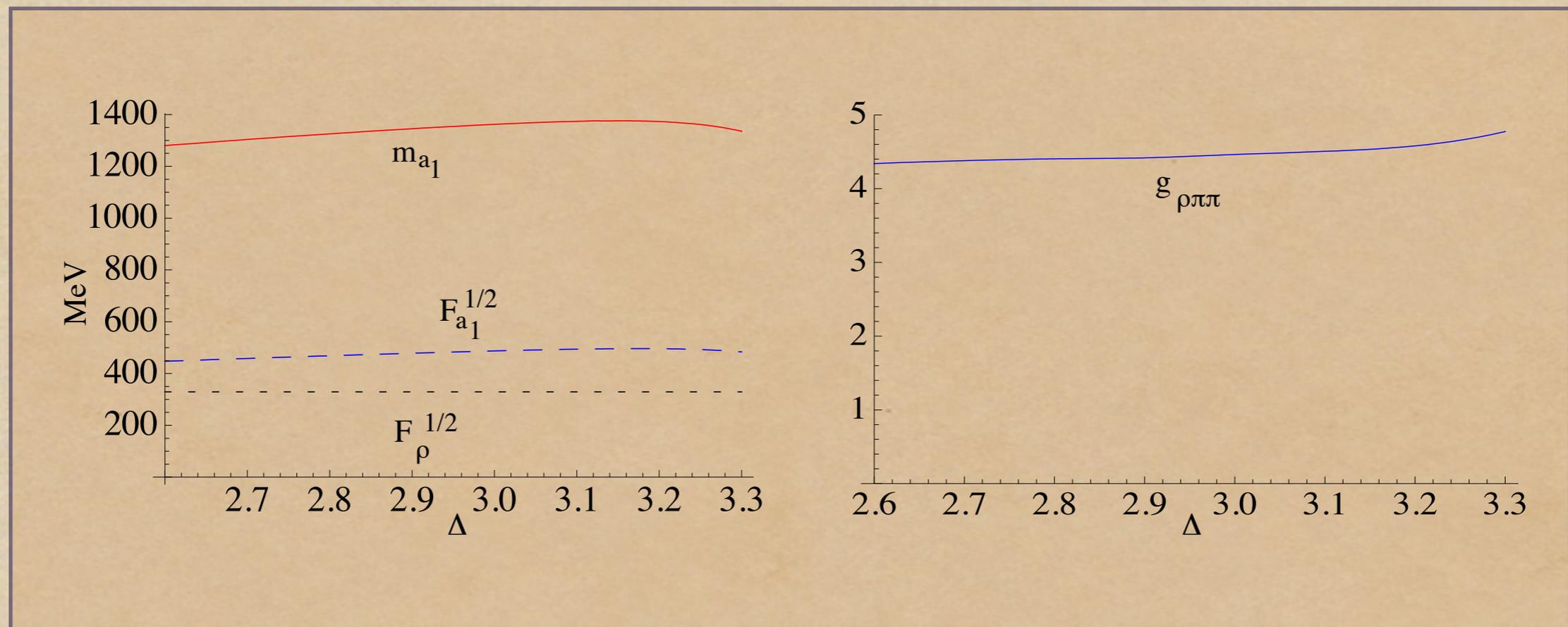
Another example: Electrical Conductivity to Charge Susceptibility σ/χ

$$\frac{\sigma}{\chi} = \frac{1}{4\pi T} \frac{d}{d-2}$$

Kovtun and Ritz '08

Universality?

Observables are roughly independent of X mass



From JE, Westenberger '09

$$m_X^2 = \Delta(\Delta - 4).$$

$\Delta = 3$ is the value set by matching to the UV.

Other Observations

- Close relationship between meson wavefunctions in extra dimension and light-front wavefunctions (Brodsky and De Teramond)
- Baryons appear as solitons in extra dimension similar to Skyrmions (Sakai and Sugimoto; Nawa *et al.*; Pomarol and Wulzer)
- AdS/QCD may address qualitative questions like chiral symmetry restoration (D.K. Hong *et al.*; Shifman and Vainshtein; Klempt)
- Can improve matching to UV by adding higher dimension 5D operators to action → power corrections in Operator Product Expansion (Hirn and Sanz)

AdS/CFT models have also been used to study the phase structure of QCD; cold atoms and superconductivity (Son; Balsubramanian *et al.*; Hartnoll *et al.*); technicolor (Hirn, Sanz; Carone *et al.*; Hong *et al.*)

Summary

- ◆ Holographic QCD models combine features of other models of QCD at low energies
- ◆ 5D tree-level calculations in holographic QCD models agree with hadronic observables at the 10-20% level, sometimes better
- ◆ 5D loop corrections and higher-dimension 5D operators have not been included -- AdS/QCD is an uncontrolled approximation above the QCD scale, so its success is a bit surprising.